

Geomagnetic data, from the Earth's surface to the core-mantle boundary

- continuous satellites records since 1999: global coverage \Rightarrow cleaner separation of internal and external sources
- observatories: continuous series, absolute intensities since 1840
- navigation: since late XVIth century, orientation only before 1840
- archeological artifacts and sediments: indirect records over the past 10,000 yrs



- Downward continuation through an electrically insulating mantle (Fig. 2) of noisy observations y^o at or above the Earth's surface
- \blacktriangleright radial potential magnetic field B_r at the core surface:

$$\mathbf{y}^o = \mathbf{H}(B_r) + \mathbf{e}^o \tag{1}$$

indirect measurements of the core state (ill-conditioned Green's functions in H)

Fig. 2- From Fournier et al (2010)





Field changes related to horizontal core motions \mathbf{u}_h through the core surface radial induction equation:

$$\frac{\partial B_r}{\partial t} = -\nabla_h \cdot (\mathbf{u}_h B_r)$$

with accuracy of the secular variation estimate depeding on both epochs and length-scales (Fig. 3)

• unresolved field features B'_r at sm $(\ell < 800 \text{ km at the core surface}),$ decadal time-scales \Rightarrow time-corre representativeness (Gillet et al, 20 $e^r = abla_h \cdot (\mathbf{u}_h)$

Fig. 3- Ensemble of realizations of the secular variation spherical harmonic coefficients dg_1^0/dt (top) and dg_{10}^5/dt (bottom), in nT/yr, with the average value in black (from Gillet et al, submitted).

Assimilating geomagnetic data into stochastic / deterministic quasi-geostrophic models of the Earth's outer core

Fig. 1- data frequency from available databases of the archeomagnetic, historical, observatory and satellite era.

$$+ \eta \nabla^2 B_r$$

mall length-scales
, with extrapolated
elated errors of
2009)
$$Br'$$

$$_{h}Br')$$
 (3)

Geomagnetic power spectrum S(f) and stochastic modeling

- Observed series suggest the process X sampled by geomagnetic records is:
- \triangleright C^1 (continuous, once differentiable) from 5 to 100 yrs periods (Fig. 6)
- spectral densities (Fig. 4) are compatible with Auto-Regressive (AR) processes and stochastic differential equations (below W stands for a white noise process)



- ▶ 100 10⁵ yrs periods: $S(f) \propto f^{-2}$ AR-1 process with $\tau_1 \sim 20,000$ yrs, e.g. $\frac{\partial X}{\partial x} + \frac{X}{\partial x} = W$
- ▶ if valid for the dipole (Fig. 5), is it still the case for smaller length-scales?

▶ 5 – 100 yrs periods: $S(f) \propto f^{-4}$

AR-2 process with $\tau_2 \sim 1000$ yrs, e.g.

$$\frac{\partial^2 X}{\partial t^2} - \frac{3X}{\tau_2^2} = W$$

 \Rightarrow the induction equation (2) suggests the flow is governed by an AR-1 stochastic differential equation of the form

$$\frac{\partial \mathbf{u}_h}{\partial t} + \frac{\mathbf{u}_h}{\tau_u} = \mathbf{W}$$

- ▶ aim at coupling (2) and (4) with an Ensemble Kalman Smoother (Evensen & Van Leeuwen, 2000), using an state augmentation approach (e.g. Reichle et al, 2002) to account for time-correlated errors in equation (3).
- ► NB: instantaneous flow accelerations are meaningless in this framework, only flow increments between two epochs make sense.

References

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 $\sim C^0$ (continuous, not differentiable) on centenial periods and longer (Fig. 5)

Fig. 5- Virtual axial dipole moment (VADM) from archeomagnetic records (Genevey et al, 2008).



(4) observatory (Germany).

This is made possible by the small size of this problem, $\mathcal{O}(500)$ parameters / epoch.

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Sampling rate of the core state

- high sampling rate in modern (observatories and satellite) data, but...
- 6 months day-time (i.e. noisy) data at high latitudes: fake global coverage
- ambiguity between core and external (magneto- and ionospheric) sources
- mantle conductivity = low-pass filter: which cut-off frequency?
- monthly to interannual periods: shall we extrapolate the AR-2 process at periods below $\mathcal{O}(1)$ year?
- should the process be different at periods shorter than the 4 yrs period of deterministic torsional (i.e. zonal and geostrophic) Alfvén waves, responsible for length-of-day (LOD) changes at 6–9 yrs periods (after Gillet et al 2010)?



Fig. 7- Torsional Alfvén waves (left), as detected from observatory records, travelling from the inner core (s = 0.35) to the equator of the outer core (s=1), and their 6–9 yrs band-pass signature (right) on the LOD changes compared to the observations.

Imaging magnetic forces inside the core with a deterministic model?

- measured by the Lundquist number

$$\mathbf{q} = \frac{1}{2H} \int_{-H}^{H} \left[B_s^2, B_\phi^2, B_s B_\phi \right] dz ,$$

 $\beta(s) \propto \frac{1}{\sqrt{ds}}$ the coriolis parameter:







band-pass 6 - 9 y ; SH degree I=9



► 3-dimensional sequential (Liu et al, 2007; Fournier et al, 2011) and variational (Li et al, 2011) attempts at estimating the core state, but filter out rapid variations ▶ in the Earth's core: inertial waves period $\tau_i \ll$ magnetic Afvén waves period $\tau_a \ll$ magnetic diffusion time τ_d \Rightarrow large length-scale transient motions invariant along the rotation axis e_z (Jault, 2008), cf. Lehnert number

$$\lambda = \tau_i / \tau_a \sim 10^{-4}$$

 \Rightarrow weak magnetic diffusion at large length-scales, as

$${m S}= au_{m d}/ au_{m a}\sim 10^5$$
 .

motivates a diffusiveless quasi-geostrophic model for z-invariant equatorial motions (following Canet et al, 2009) Fig. 8- Example of equatorial maps for ψ (top)

 $\mathbf{u}_{e}(s,\phi,t) = \mathbf{e}_{z} \times \nabla \psi$, and the *z*-averaged quadratic quantities

with H(s) the half-height of a fluid column, and

$$\beta \frac{\partial^2 \psi}{\partial t \partial \phi} = \mathcal{F}(\psi, \mathbf{q})$$

coupling eqns (1), (2) and (5) using an EnKF ► to obtain a first image of magnetic forces **q**: to forecast the field evolution (candidate to the International Geomagnetic Reference Field)

(a) must account for large uncertainties on ψ , even more on $\partial_t \psi$ (b) favored by a suspected steep spatial power spectrum for **q** (c) bounded value problem (positiveness + Cauchy-Schwartz constraint on the unknown \mathbf{q}) (d) reconcile stochastic constraint (4) with (5): consider flow increments instead of $\partial_t \psi$.



and $d\psi/dt$ (bottom) obtained from satellite data in 2005 and equation (2), after Gillet et al (2009).

